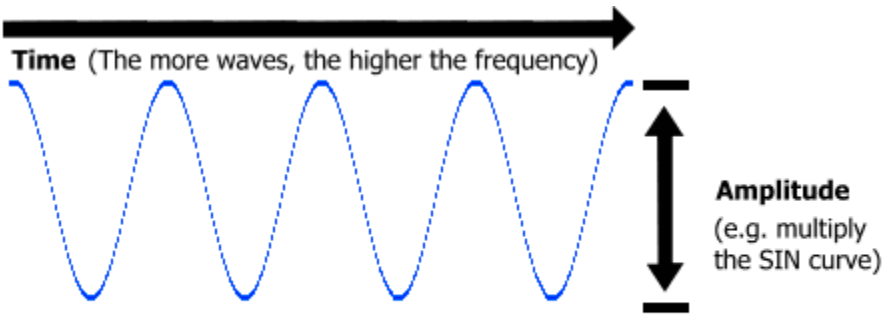


# **A Plausibility Argument for the Uncertainty Principle**

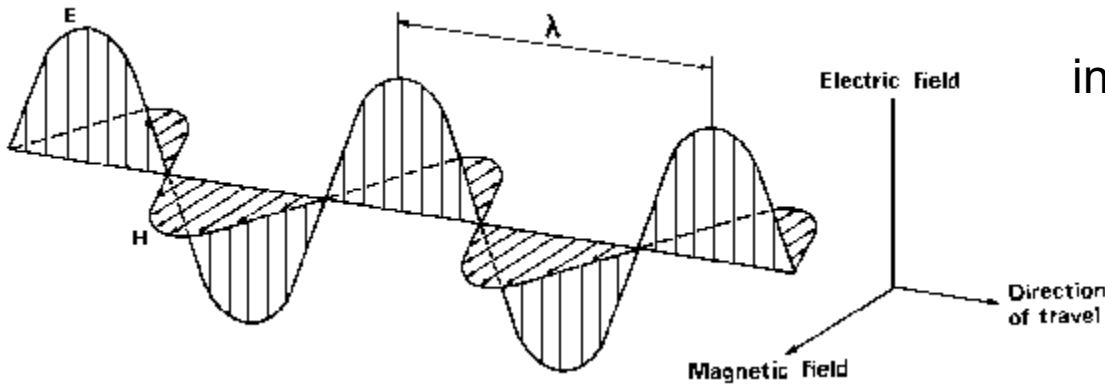
# waves go on forever...



A wave that propagates forever in one dimension is described by:

$$E = E_0 \cos\left(\frac{2\pi x}{\lambda} - 2\pi f t\right)$$

$$c = \lambda f$$



in shorthand:

$$E = E_0 \cos(kx - \omega t)$$

$$\omega = 2\pi f, \quad k = 2\pi / \lambda$$

angular frequency      wave number

In the quantum description of a particle, the wavelength represents the momentum

$$\lambda_{dB} = \frac{h}{p}$$

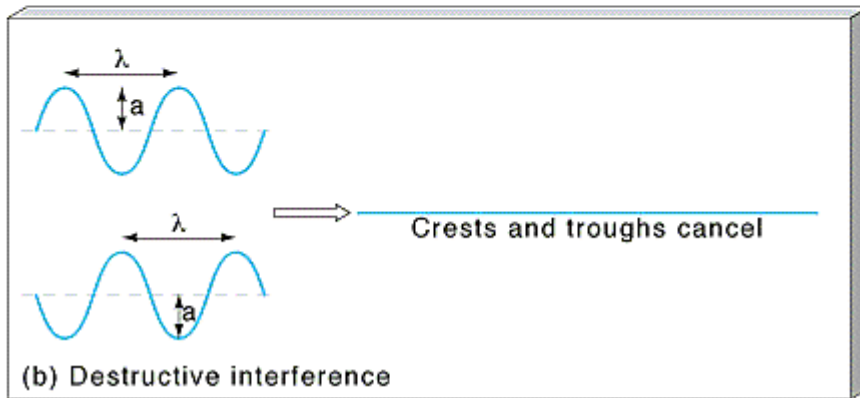
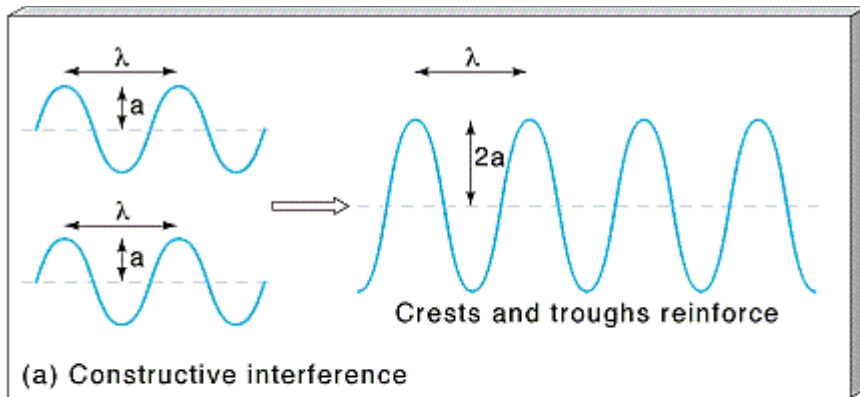
$$\Psi(x, t) = \Psi_0 e^{i(kx - \omega t)}$$

For this kind of wave the momentum has a single value, with no uncertainty

$$p = \hbar k \quad \text{and} \quad \sigma_p = 0 \quad \text{or} \quad \Delta p = 0$$

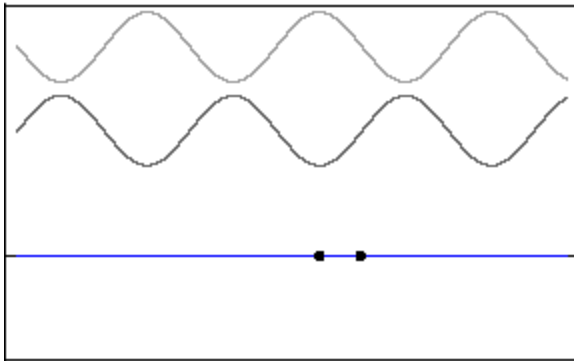
But \*where\* is the wave????

# Interference



**Because of linear superposition waves can interfere (add or cancel)**

**Recall that we built linearity into the Schrodinger equation!**



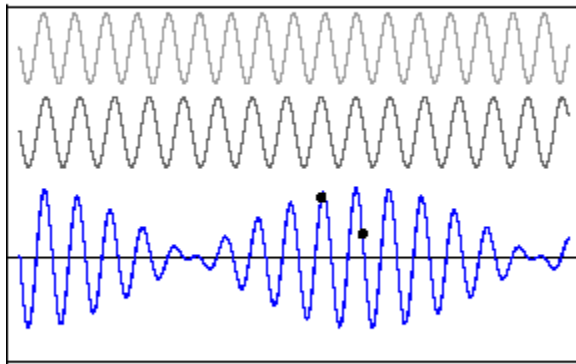
# How to Make a Localized Wave?

Interfering waves, generally...

$$y = y_1 + y_2 = A \cos(k_1 x - \omega_1 t) + A \cos(k_2 x - \omega_2 t)$$

⇓

$$y = 2A \cos \frac{1}{2} \{ (k_2 - k_1)x - (\omega_2 - \omega_1)t \} \bullet \cos \frac{1}{2} \{ (k_1 + k_2)x - (\omega_1 + \omega_2)t \}$$



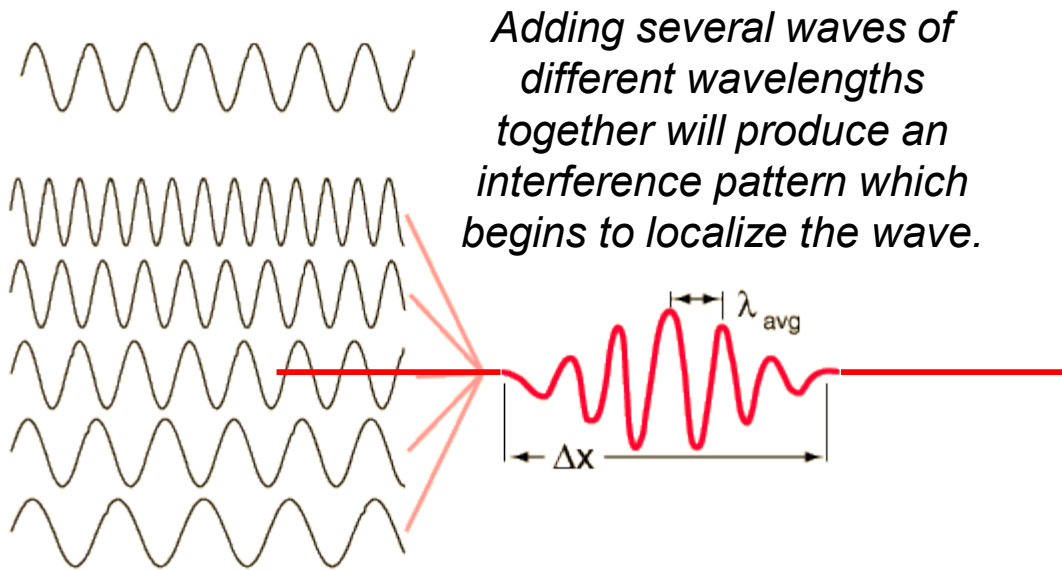
“Beats” occur when you add two waves of slightly different frequency. They will interfere constructively in some areas and destructively in others.

Can be interpreted as a sinusoidal envelope:

$$2A \cos \left( \frac{\Delta k}{2} x - \frac{\Delta \omega}{2} t \right)$$

Modulating a high frequency wave within the envelope:  $\cos \left[ \frac{1}{2} (k_1 + k_2)x - \frac{1}{2} (\omega_1 + \omega_2)t \right]$

**FOURIER THEOREM:** any wave packet can be expressed as a superposition of an infinite number of harmonic waves



To form a pulse that is zero everywhere outside of a finite spatial range  $\Delta x$  requires adding together an infinite number of waves with continuously varying wavelengths and amplitudes.

spatially localized wave group

adding varying amounts of an infinite number of waves

sinusoidal expression for harmonics

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} a(k) e^{ikx} dk$$

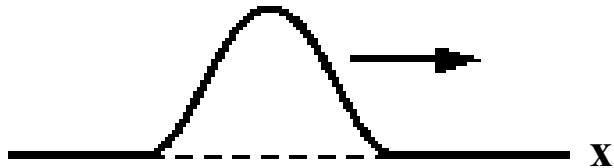
amplitude of wave with wavenumber  $k=2\pi/\lambda$

# *The Uncertainty Principle*

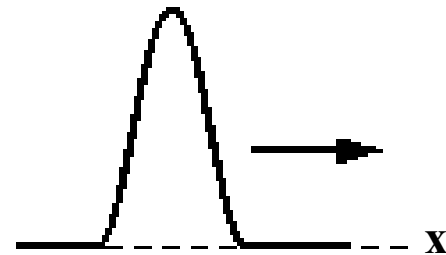
**Remember the cosine wave that went on “forever”?**

We knew its momentum very precisely, because the momentum is a function of the wavelength, and the wavelength was very well defined.

But what is the wavelength of a localized wave packet? We had to add a bunch of waves of different wavelengths to produce it.



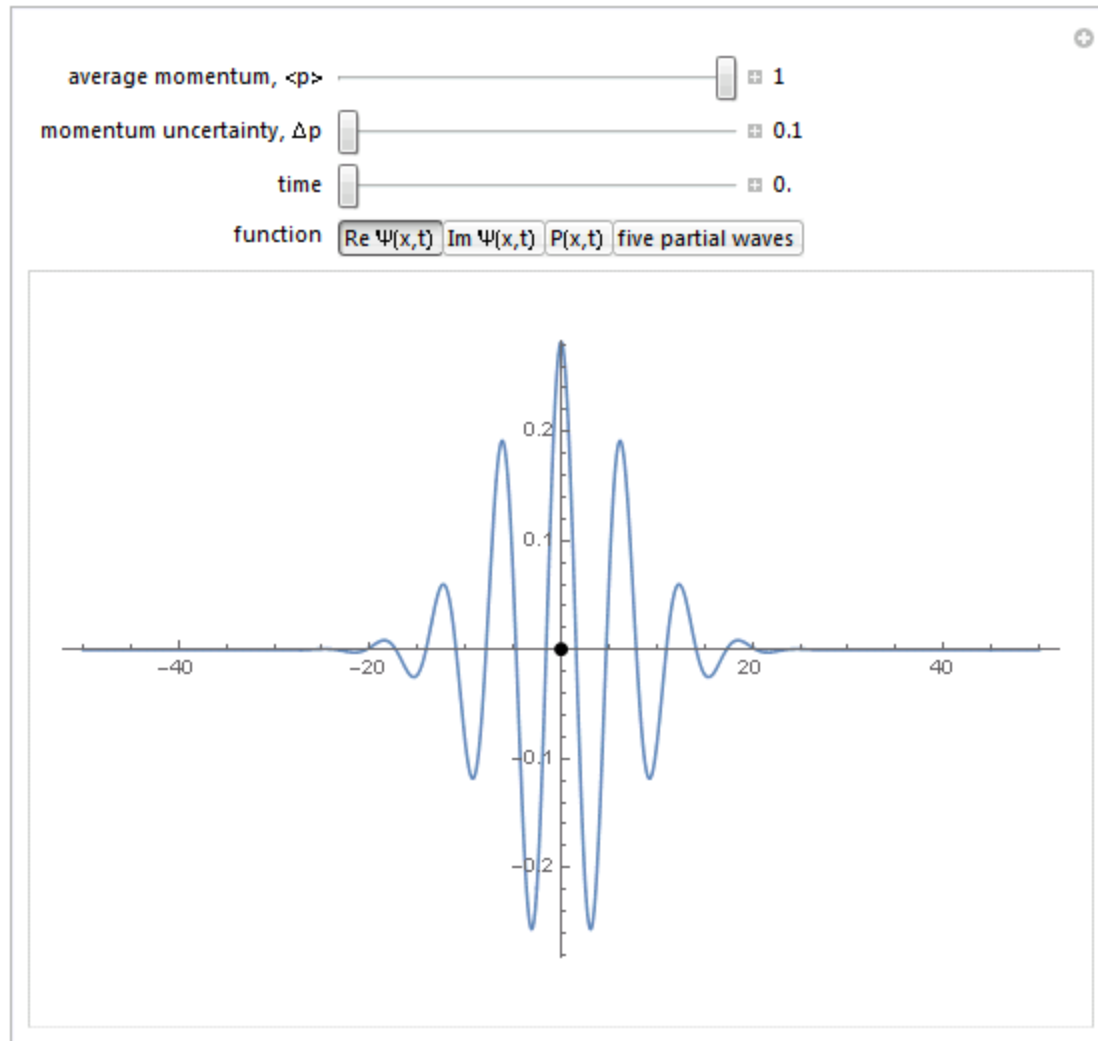
Producing a localized wave packet requires a continuous distribution of wavelengths.



Producing a wave packet which is more confined in space requires a wider distribution of wavelengths.

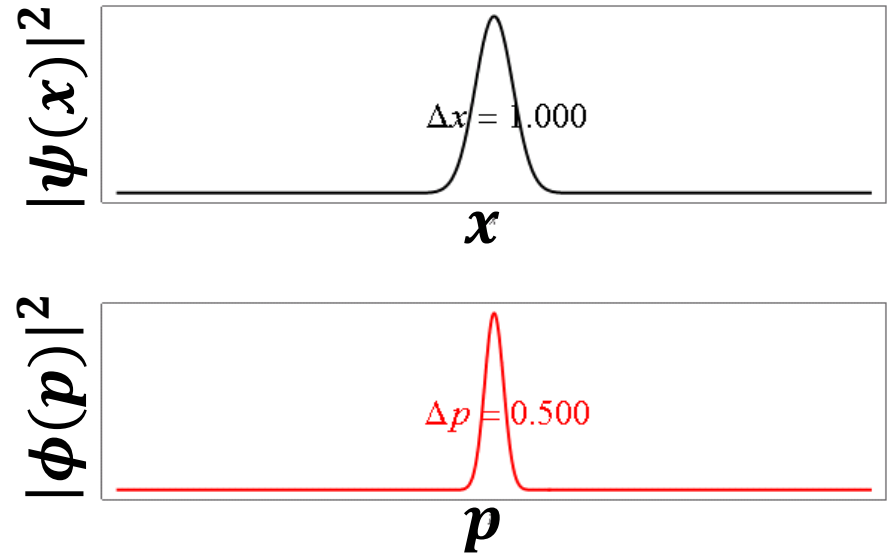
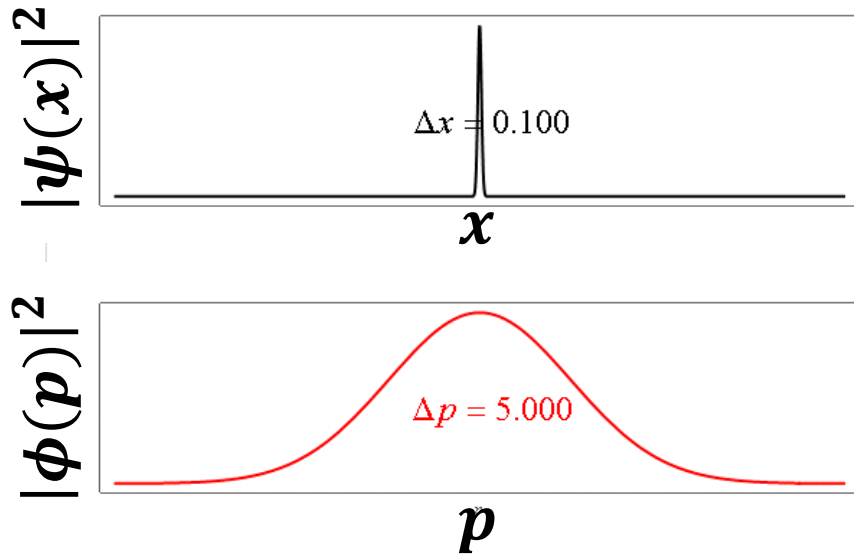
**Consequence:** The more localized the wave packet in space, the less precisely defined is the momentum.

## Wavepacket for a Free Particle



<http://demonstrations.wolfram.com/WavepacketForAFreeParticle/>

# Uncertainty Relation for Gaussian Wavepackets

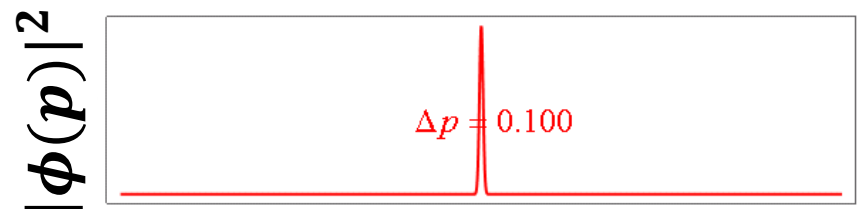
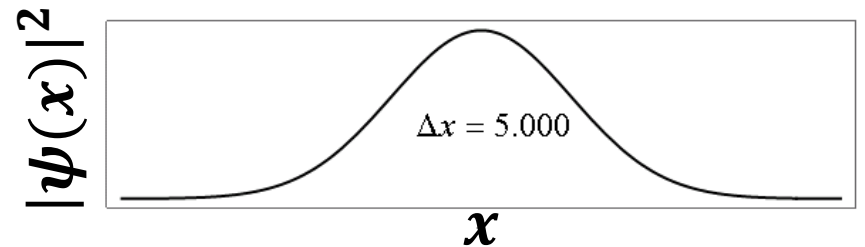


$$\Delta x \Delta p = 1/2$$

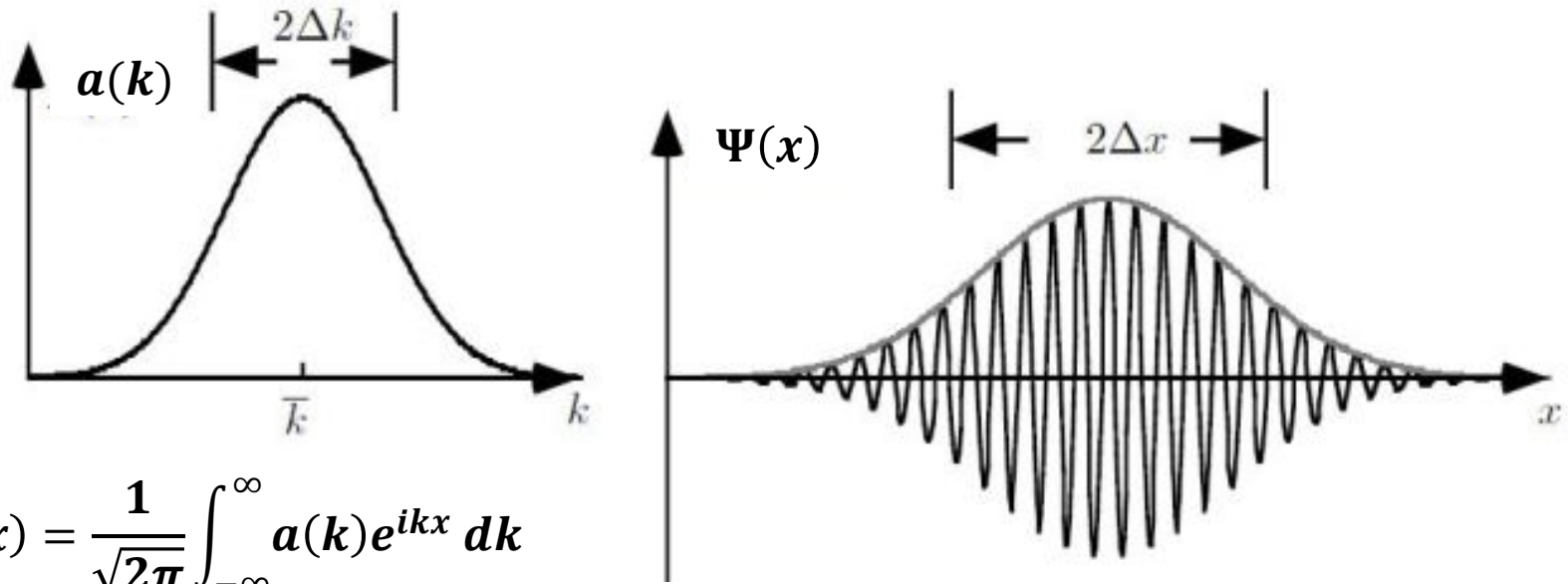
( $\hbar = 1$  in this simulation)

Later we will show:

$$\sigma_x \sigma_p \geq \hbar/2$$







$$\Psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} a(k) e^{ikx} dk$$

## Wave Uncertainty Relation

$$\Delta x \cdot \Delta k \sim 1$$

Translating to Quantum Mechanics:

$$\Delta x \cdot \hbar \Delta k \sim \hbar \quad \text{or} \quad \Delta x \cdot \Delta p_x \sim \hbar$$

Ultimately we will show that  $\sigma_x \cdot \sigma_{p_x} \geq \hbar/2$